

An Uncertainty Principle like Hardy's Theorem for Nilpotent Lie Group G_n

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Abstract

Let $f: \mathfrak{R} \rightarrow \mathbb{C}$ be measurable and for all $x, y \in \mathfrak{R}$ and if

$$(i) \quad |f(x)| \leq C \exp(-a\pi x^2)$$

$$(ii) \quad |\hat{f}(y)| \leq C \exp(-b\pi y^2)$$

where $C, a, b > 0$. If $ab > 1$ then $f = 0$ a.e. If $ab = 1$ then $f(x) = C \exp(-a\pi x^2)$. If $\alpha\beta < 1$ then there exist infinitely many linearly independent function satisfying (i) and (ii).

In this paper we extend an uncertainty principle due to Cowling and Price to threadlike nilpotent Lie groups G_n .

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Introduction

A classical theorem of Hardy [6] on Fourier transform pairs says that a non-zero function f on the real line \mathfrak{R} and its Fourier transform \hat{f} can not both be very rapidly decreasing. More precisely, let the Fourier transform be defined by,

$$\hat{f}(y) = \int_{\mathfrak{R}} f(x) \exp(-2\pi ixy) dx, y \in \mathfrak{R}$$

The following is a generalization of this theorem due to Cowling and Price [3].

Theorem (Cowling and Price): Let $f: \mathfrak{R} \rightarrow \mathbb{C}$ be measurable and

$$(i) \quad \|e_a f\|_{L^p(\mathfrak{R})} < \infty$$

$$(ii) \quad \|e_b \hat{f}\|_{L^q(\mathfrak{R})} < \infty$$

Where $a, b > 0$, $e_k(x) = \exp(k\pi x^2)$ and $1 \leq \min(p, q) < \infty$. If $ab \geq 1$, then $f = 0$ a.e. If $ab < 1$, then there exist infinitely many linearly independent functions satisfying (i) and (ii)

An analogue of the Cowling-Price theorem has been proved in [3] for Euclidian space the Heisenberg group H_n and the Euclidian motion group of the plane.

Main Results

Theorem: let $f: G_n \rightarrow \mathbb{C}$ be a measurable function such that $|f(x)| \leq C \exp(-a\pi \|x\|^2)$ for some $c, a > 0$ and all $x \in G_n$ then the function $h(\xi) = \|\pi_\xi(f)\|_{HS}$ is bounded.

Proof:

$$\begin{aligned} |\xi_1| \|\pi_\xi(f)\|_{HS}^2 &= \int_{\mathfrak{R}^2} |\mathcal{F}_1 \dots (n-1) f(\xi_1, t, q_3(\xi, t), q_{n-1}(\xi, t), s)|^2 dt ds \\ &\leq \int_{\mathfrak{R}^{n-1}} |f(x_1, \dots, x_{n-1}, s)| |e^{2\pi i \xi_1 x_1}| |e^{2\pi i t x_2}| dx_1 dx_2 \dots dx_{n-1} \\ &\leq C \int_{\mathfrak{R}^{n-1}} \exp(-a\pi(x_1^2 + \dots + x_n^2 + s^2)) |\exp(2\pi i \xi_1 x_1)| |\exp(2\pi i x_2)| dx_1 \dots dx_{n-1} \\ &\leq \text{constant} \exp(-a\pi s^2) \int_{\mathfrak{R}^2} |\exp(-a\pi(x_1^2 - \frac{2}{a} i \xi_1 x_1))| |\exp(-a\pi(x_2^2 - \frac{2}{a} i t x_2))| dx_1 dx_2 \end{aligned}$$

$$\int_{-\infty}^{\infty} |\exp(-a\pi(x^2 - \frac{2}{a} i \xi_1 x))| dx = \int_{-\infty}^{\infty} |\exp(-a\pi(x + iy)^2 - \frac{2}{a} i \xi_1(x + iy))| dx$$

By the change of contour $x \rightarrow x + iy$

$$\begin{aligned} &= \int_{-\infty}^{\infty} \exp \operatorname{Re}(-a\pi((x + iy)^2 - \frac{2}{a} i \xi_1(x + iy))) dx \\ &= e^{-2\pi \xi_1 y} \int_{-\infty}^{\infty} \exp(-a\pi(x^2 - y^2)) dx \\ &= e^{-2\pi(\xi_1 y - \frac{a}{2} y^2)} \int_{-\infty}^{\infty} \exp(-a\pi x^2) dx \\ &\leq \text{const. } e^{-2\pi(\xi_1 y - \frac{a}{2} y^2)} \end{aligned}$$

Taking infimum over y , we have

$$\begin{aligned} \int_{-\infty}^{\infty} |\exp(-a\pi(x^2 - \frac{2}{a}i\xi_1 x))| dx &\leq \text{const } e^{-2\pi} \sup \left(\xi_1 y - \frac{a}{2} y^2 \right) \\ &= \text{const } e^{-2\pi\xi_1^2/2a} \\ &= \text{const } e^{-\pi\xi_1^2/a} \end{aligned}$$

$$\therefore \|\pi_{\xi}(f)\|_{HS} \leq \text{const } (|\xi_1|^{-1} \exp(-\pi\xi_1^2/a))^{1/2}$$

Thus h is bounded.

The sharpness of the constant

Let G_4 be the low dimensional nilpotent Lie groups for $ab < 1$, let $\alpha \in (a, 1/b)$, but it works for G_n using lemmas 2.1, 2.2 in [7]

For $x \in G_4$, let $f(x) = \exp(-\alpha\pi \|x\|^2)$

We know that,

$$\|\pi_{\xi_1, \xi_2}(f)\|_{HS}^2 = |\xi_1|^{-1} \int_{\mathfrak{R}} |\mathcal{F}_{123} f(\xi_1, u, \frac{1}{2} \frac{u^2}{\xi_1} + \xi_3, v)|^2 du dv$$

For $u, v \in \mathfrak{R}$,

$$\begin{aligned} \mathcal{F}_{123} f(\xi_1, u, \frac{1}{2} \frac{u^2}{\xi_1} + \xi_3, v) &= \int_{\mathfrak{R}^3} e^{-\pi(x_1^2 + x_2^2 + x_3^2 + v^2)} e^{-2\pi i x_1 \xi_1} e^{-2\pi i x_2 u} e^{-2\pi i x_3 (\xi_3 + (1/2)(u^2/\xi_1))} dx_1 dx_2 dx_3 \\ &= e^{-\pi v^2} \mathcal{F}(e^{-\pi x^2 1^2})(\xi_1) \mathcal{F}(e^{-\pi x 2^2})(u) \mathcal{F}(e^{-\pi x 3^2})(\xi_3 + (1/2)(u^2/\xi_1)) \dots (A) \\ &= e^{-\pi v^2} \mathcal{F}(e^{-\pi x 1^2})(\xi_1) \mathcal{F}(e^{-\pi x 2^2})(u) \mathcal{F}(e^{-\pi x 3^2 - 2\pi i x_3 x_3} ((1/2)(u^2/\xi_1))) \end{aligned}$$

where $\mathcal{F}(e^{-\pi x 1^2})(\xi_1)$ is the Fourier transform of the function $g(x) = e^{-\pi x^2}$ at ξ_1

Now, $\mathcal{F}(e^{-\pi x 1^2})(\xi_1) = \text{const } e^{-\pi \xi_1^2/4}$

And $\mathcal{F}(e^{-\pi x 2^2})(u) = \text{const } e^{-\pi u^2/4}$

So, $|\mathcal{F}_{123} f(\xi_1, u, \frac{1}{2} \frac{u^2}{\xi_1} + \xi_3, v)|$

$$= \text{const. } e^{(-\pi \xi_1^2)/4} e^{(-\pi u^2)/4} e^{-\pi v^2} |\mathcal{F}(e^{-2\pi i \xi_3 x_3 - \pi x 3^2}) \left(\frac{1}{2} \frac{u^2}{\xi_1} \right)|$$

$$\leq \text{const. } e^{(-\pi \xi_1^2)/4} e^{(-\pi u^2)/4} e^{-\pi v^2} \int_{\mathfrak{R}} |e^{-2\pi i \xi_3 x_3 - \pi x 3^2}| dx_3$$

$$\begin{aligned}
&= \text{const. } e^{(-\pi\xi_1^2)/\alpha} e^{(-\pi u^2)/\alpha} e^{-\pi uv^2} e^{(-\pi\xi_3^2)/\alpha} \int_{\mathfrak{R}} |e^{-\alpha\pi\{\xi_3 + (i\xi_3/\alpha)^2\}^2}| dx_3 \\
&= \text{const. } e^{(-\pi\xi_1^2)/\alpha} e^{(-\pi u^2)/\alpha} e^{-\pi uv^2} e^{(-\pi\xi_3^2)/\alpha} \\
\therefore \|\pi_{\xi_1, \xi_3}(f)\|_{\text{HS}}^2 &\leq \text{const. } |\xi_1|^{-1} e^{-2\pi(\xi_1^2 + \xi_3^2)/\alpha} \int_{\mathfrak{R}} e^{-2\pi u^2/\alpha} e^{-2\pi v^2} du dv \\
&= \text{const. } |\xi_1|^2 e^{-2\pi(\xi_1^2 + \xi_3^2)/\alpha} \dots \text{(B)}
\end{aligned}$$

$$\begin{aligned}
&\int_{\mathfrak{R}^2} e^{2\pi b(\xi_1^2 + \xi_3^2)} |\xi_1| \|\pi_{\xi_1, \xi_3}(f)\|_{\text{HS}}^2 d\xi_1 d\xi_3 \\
&\leq \text{const. } \int_{\mathfrak{R}^2} e^{2\pi(b-1/\alpha)(\xi_1^2 + \xi_3^2)} d\xi_1 d\xi_3 \\
&< \infty \text{ since } b - \frac{1}{\alpha} < 0
\end{aligned}$$

Thus there is a non zero function satisfying the condition (i) and (ii) of theorem 2.4 [7] for $ab < 1$.

Since $|f(x)| = \exp(-\alpha\pi x^2) < \exp(-\pi x^2)$ for $q = 2$.

In case the condition (ii) is replaced by

$$\int_{\mathfrak{R}^2} e^{2\pi b\xi_1^2} |\xi_1| \|\pi_{\xi_1, \xi_3}(f)\|_{\text{HS}}^2 d\xi_1 d\xi_3 < \infty$$

which is really the condition used in the proof of $q = 2$, the things are more simpler.

Write the equation (A) as

$$\begin{aligned}
&\mathcal{F}_{123} f\left(\xi_1, u, \frac{1}{2} \frac{u^2}{\xi_1} + \xi_3, v\right) \\
&= \text{const. } e^{-\alpha\pi v^2} e^{(-\pi\xi_1^2)/\alpha} e^{(-\pi u^2)/\alpha} e^{-(\pi/\alpha)\{\xi_3 + (1/2)(u^2/\xi_1)\}^2} \\
\therefore \|\pi_{\xi_1, \xi_3}(f)\|_{\text{HS}}^2 &= \text{const. } |\xi_1|^{-1} e^{(-2\pi\xi_1^2)/\alpha} \int_{\mathfrak{R}^2} e^{-2\alpha\pi v^2} e^{-2\pi/\alpha\{u^2 + \xi_3 + (1/2)(u^2/\xi_1)\}^2} du dv \\
&= \text{const. } |\xi_1|^{-1} e^{(-2\pi\xi_1^2)/\alpha} \int_{\mathfrak{R}} e^{-2\pi/\alpha\{u^2 + (\xi_3 + 1/2 u^2/\xi_1)\}^2} du \\
&\int_{\mathfrak{R}^2} e^{-2\pi b\xi_1^2} |\xi_1| \|\pi_{\xi_1, \xi_3}(f)\|_{\text{HS}}^2 d\xi_1 d\xi_3 \\
&= \text{const. } \int_{\mathfrak{R}^2} e^{2\pi(b-1/\alpha)\xi_1^2} e^{(-2\pi u^2)/\alpha} e^{-2\pi/\alpha\{\xi_3 + (1/2)(u^2/\xi_1)\}^2} du d\xi_1 d\xi_3
\end{aligned}$$

Applying $\xi_3 \rightarrow \xi_3 - \frac{1}{2} \frac{u^2}{\xi_1}$

$$= \text{const.} \int_{\mathfrak{R}^3} e^{2\pi\{b-1/\alpha\} \xi_1^2} e^{(-2\pi u^2)/\alpha} e^{-2\pi/\alpha \xi_3^2} du d\xi_1 d\xi_3 < \infty$$

For $q \geq 2$ or $1 \leq q < 2$ and $ab < 1$.

$$\begin{aligned} & \int_{\mathfrak{R}^2} e^{qb\pi\{\xi_1^2 + \xi_3^2\}} |\xi_1|^{|\pi\xi} \xi_3 (f) \|_{HS}^q d\xi_1 d\xi_3 \dots(1) \\ & \int_{\mathfrak{R}^2} e^{qb\pi\{\xi_1^2 + \xi_3^2\}} |\xi_1|^{\{1-(q/2)\}} e^{-(ap/\alpha)(\xi_1^2 + \xi_3^2)} d\xi_1 d\xi_3 \\ & = \int_{\mathfrak{R}^2} e^{\pi q(b-1/\alpha)(\xi_1^2 + \xi_3^2)} |\xi_1|^{1-q/2} dx_1 d\xi_3 < \infty \text{ for } 1 \leq q < 2 \end{aligned}$$

and for $q \geq 2$ considering the case $|\xi_1| > 1$ we obtain integral in (i) is finite.

Theorem: Let a and b be positive real numbers and $1 \leq \min(p, q) < \infty$. Suppose that

$f \in L^1(G_N) \cap L^2(G_N)$ satisfies the following conditions:

- (i) $\int_{G_N} e^{pa\pi \|\xi\|^2} |f(x)|^p dx < \infty$
- (ii) $\int_{\mathfrak{R}^{n-2}} |\xi_1| e^{b\pi q \|\xi\|^2} \|\pi_\xi(f)\|_{HS}^q d\xi < \infty$

If $1 \leq q < 2$ and $ab \geq \frac{2}{q}$ then $f = 0$ a.e.

Proof: Let p be such that $\frac{1}{p} + \frac{1}{q} = 1$. Clearly $p \geq 2$ and $q < p$ by Lemma 2.2 [7] we have

$$\begin{aligned} \|\mathbf{e}_{b\{1+(q/p)\}} \hat{g}\|_1 &= \int_{\mathfrak{R}} \mathbf{e}_{b\{1+(q/p)\}}(\xi_1) |\hat{g}(\xi_1)| d\xi_1 \\ &= \int_{\mathfrak{R}^{n-2}} |\xi_1| \mathbf{e}_{b\{1+(q/p)\}}(\xi_1) \|\pi_\xi(f)\|^2 d\xi_1 d\xi_3 \dots d\xi_{n-1} \end{aligned}$$

Define the function u and v in \mathfrak{R}^{n-2} by

$$\begin{aligned} u(\xi) &= \mathbf{e}_b(\xi_1) |\xi_1|^{1/q} \|\pi_\xi(f)\|_{HS} \text{ and} \\ v(\xi) &= \mathbf{e}_{bq/p}(\xi_1) |\xi_1|^{1/p} \|\pi_\xi(f)\|_{HS} \end{aligned}$$

Then,

$$\int_{\mathfrak{R}^{n-2}} |u(\xi)|^q d\xi = \int_{\mathfrak{R}^{n-2}} \mathbf{e}_{bq}(\xi_1) |\xi_1| \|\pi_\xi(f)\|_{HS}^q d\xi < \infty$$

and

$$\int_{\mathfrak{R}^{n-2}} |v(\xi)|^p d\xi = \int_{\mathfrak{R}^{n-2}} \mathbf{e}_{bq}(\xi_1) |\xi_1| \|\pi_\xi(f)\|_{HS}^p d\xi$$

$$\begin{aligned}
&= \int_{\mathfrak{R}^{n-2}} e_{bq}(\xi_1) |\xi_1| \|\pi_{\xi}(f)\|_{HS}^q \|\pi_{\xi}(f)\|_{HS}^{p-q} d\xi \quad (p > q) \\
&\leq K \int_{\mathfrak{R}^{n-2}} e_{bq}(\xi_1) |\xi_1| \|\pi_{\xi}(f)\|_{HS}^q > \infty
\end{aligned}$$

So using Holders inequality we have

$$\|e_{b(1+(q,p))} \hat{g}\|_1 \leq \|u\|_q \|v\|_p < \infty$$

$$\text{Since } \frac{a}{2} b \left(1 + \frac{q}{p}\right) \geq \frac{1}{2} \cdot \frac{2}{q} \left(1 + \frac{q}{p}\right) = \frac{1}{2} \cdot \frac{2}{q} \cdot q = 1$$

The Cowling price theorem shows that $g = 0$ and hence $f = 0$ a.e.

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